## **Topology of Scalar and Vector Fields**

# G.A. GIRALDI<sup>1</sup>

<sup>1</sup>LNCC-National Laboratory for Scientific Computing -Av. Getulio Vargas, 333, 25651-070, Petrópolis, RJ, Brazil {gilson}@lncc.br

Abstract. .

#### **1** Introduction

In this Chapter, we review mathematical elements in scalar and vector fields. We start with basic elements in linear algebra (see [1, 2] for details in this area). Next, in section 3-4 we define scalar and vector fields in  $\mathbb{R}^n$ .

### 2 Background in Linear Algebra

Let the N-dimensional vector space V over  $\Re$  composed by n-uplas:

$$\mathbf{v} = \begin{pmatrix} v (0) \\ v (1) \\ \vdots \\ \vdots \\ v (N-1) \end{pmatrix}.$$
 (1)

and B a basis of V given by N linearly independent vectors:

$$B = \{ \mathbf{e}_i \in V; \quad , \quad i = 0, 1, ..., N - 1 \}.$$
(2)

Therefore, we know that any vector  $\mathbf{v} \in V$  can be written as a linear combination of elements in B, that means:

$$\mathbf{v} = \sum_{i=0}^{N-1} \alpha_i \mathbf{e}_i. \tag{3}$$

We call the array composed by the coefficients  $\alpha_i$  the representation of v in the basis B, which we indicate by:

$$[\mathbf{v}]_B = \begin{pmatrix} \alpha (0) \\ \alpha (1) \\ \vdots \\ \vdots \\ \alpha (N-1) \end{pmatrix}.$$
(4)

Moreover, we know that such representation is unique. If B is the canonical basis given by:

$$B = \left\{ \begin{pmatrix} 1 \\ 0 \\ \cdot \\ \cdot \\ \cdot \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 1 \\ \cdot \\ \cdot \\ \cdot \\ 0 \end{pmatrix}, \begin{pmatrix} 0 \\ 0 \\ 1 \\ \cdot \\ \cdot \\ 0 \end{pmatrix}, \cdots, \begin{pmatrix} 0 \\ 0 \\ \cdot \\ \cdot \\ \cdot \\ 1 \end{pmatrix} \right\},$$
(5)

them we have  $[\mathbf{v}]_B \equiv \mathbf{v}$ . In this case, we will drop the subscript "B" in expression (4).

In this context, we can consider functions  $T: V \to \Re$ , that satisfy the property:

$$T\left(\alpha_{1}\mathbf{u}_{1}+\alpha_{2}\mathbf{u}_{2}\right)=\alpha_{1}T\left(\mathbf{u}_{1}\right)+\alpha_{2}T\left(\mathbf{u}_{2}\right),\tag{6}$$

 $\alpha_1, \alpha_2 \in \Re$  and  $\mathbf{u}_1, \mathbf{u}_2 \in V$ . The linear function T is called a *linear form*.

The space of linear forms is a vector space over  $\Re$ , called the dual of V, which is denoted by  $V^*$ . Let the dual basis

$$B = \left\{ \mathbf{e}^{i} \in V^{*}; \quad , \quad i = 0, 1, ..., N - 1 \right\}.$$
(7)

where  $e^i$  is the dual of the vector  $e_i$ ; defined by:

$$\mathbf{e}^{i}\left(\mathbf{e}_{j}\right) = \delta\left(i-j\right)$$

Observe that  $(V^*)^* \equiv V$ . Therefore, we can see  $\mathbf{u} \in \mathbf{V}$  as a linear form  $\mathbf{u} : V^* \to \Re$ ;

$$\mathbf{u}\left(T\right) = T\left(\mathbf{u}\right)$$

### **3** Scalar and Vector Fields in $\mathbb{R}^n$

The function

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is a scalar field in  $\mathbb{R}^n$ .

If V is vector space, the function

 $\boldsymbol{F}:\mathbb{R}^n\to V,$ 

 $f: \mathbb{R}^n \to \mathbb{R},$ 

is a vector field from  $\mathbb{R}^n$  to the vector space V. A specific and important case of vector field is:

$$\boldsymbol{F}:\mathbb{R}^3\to\mathbb{R}^3,$$

# 4 Topology of Scalar Fields

Given a scalar field  $f : \mathbb{R}^n \to \mathbb{R}$  we say that a point  $\mathbf{x}_0 \in \mathbb{R}^n$  is a critical point of f if  $D_{\mathbf{x}} f(\mathbf{x}_0) = \mathbf{0}$ ; that means:

$$\frac{\partial f}{\partial x_1}(\mathbf{x}_0) = \frac{\partial f}{\partial x_2}(\mathbf{x}_0) = \dots = \frac{\partial f}{\partial x_n}(\mathbf{x}_0) = 0.$$
(8)

Also, we can compute the Hessian  $H_{\mathbf{x}}f(\mathbf{x}_0)$  as:

$$\left[H_{\mathbf{x}}f\left(\mathbf{x}_{0}\right)\right]_{ij} = \frac{\partial^{2}f}{\partial x_{i}\partial x_{j}}\left(\mathbf{x}_{0}\right), \quad 1 \le i, j \le n,$$

$$\tag{9}$$

and the dimension of the null space of the matrix  $[H_{\mathbf{x}}f(\mathbf{x}_0)]$  is called the corank of  $H_{\mathbf{x}}f$  at  $\mathbf{x}_0 \in \mathbb{R}^n$ .

We say that f has a nondegenerate critical point at  $\mathbf{x}_0 \in \mathbb{R}^n$  if  $D_{\mathbf{x}}f(\mathbf{x}_0) = \mathbf{0}$  and if  $H_{\mathbf{x}}f(\mathbf{x}_0)$ , defined by expression (9), is non-singular (nondegenerate quadratic form). On the other hand, if  $H_{\mathbf{x}}f(\mathbf{x}_0)$  is degenerate we say the point  $\mathbf{x}_0 \in \mathbb{R}^n$  that satisfies  $D_{\mathbf{x}}f(\mathbf{x}_0) = \mathbf{0}$  is a degenerate critical point.

Taylor Series for a crtical point  $\mathbf{x}_0 \in \mathbb{R}^n$ :

$$f(\boldsymbol{x}) = f(\mathbf{x}_0) + \frac{1}{2!} (\boldsymbol{x} - \mathbf{x}_0)^T H_{\mathbf{x}} f(\mathbf{x}_0) (\boldsymbol{x} - \mathbf{x}_0) + O(||\boldsymbol{x} - \mathbf{x}_0||^2).$$
(10)

From expression (10) we have tree possibilities for the critical point  $\mathbf{x}_0 \in \mathbb{R}^n$ :

- Local maximum:  $D_{\mathbf{x}}f(\mathbf{x}_0) = \mathbf{0}$  and  $H_{\mathbf{x}}f(\mathbf{x}_0)$  positive definite
- Local minimum:  $D_{\mathbf{x}}f(\mathbf{x}_0) = \mathbf{0}$  and  $H_{\mathbf{x}}f(\mathbf{x}_0)$  negative definite
- Seddle point:  $D_{\mathbf{x}}f(\mathbf{x}_0) = \mathbf{0}$  and  $H_{\mathbf{x}}f(\mathbf{x}_0)$  non-singular but with positive and negative eigenvalues
- Degenerate point:  $D_{\mathbf{x}}f(\mathbf{x}_0) = \mathbf{0}$  and  $H_{\mathbf{x}}f(\mathbf{x}_0)$  singular

## 5 Topology of Vector Fields

Consider the vector field:

$$F: \mathbb{R}^n \to \mathbb{R}^n,$$

where:

$$\boldsymbol{F}\left(\boldsymbol{x}\right)=A\boldsymbol{x},$$

with A a real and constant matrix.

In this case, the solution of the differential equation:

$$\frac{d\boldsymbol{x}}{dt} = A\boldsymbol{x},\tag{11}$$

subject to:

$$\boldsymbol{x}\left(0\right) = \boldsymbol{x}_{0},\tag{12}$$

is [3]:

$$\boldsymbol{x}(t) = \exp\left(At\right)\boldsymbol{x}_{0}.$$

We shall remember that:

$$\exp(At) = I + At + \frac{1}{2!} (At)^2 + \frac{1}{3!} (At)^3 + \dots$$

Therefore, if  $(v_i, \lambda_i)$ , i = 1, 2, ..., n are the distint eigenvectors/eigenvalues of matrix A, then, we can write:

$$\boldsymbol{x}_0 = \sum_{i=1}^n \alpha_i \boldsymbol{v}_i,\tag{13}$$

and:

$$\boldsymbol{x}(t) = \exp(At) \, \boldsymbol{x}_0 = \exp(At) \left[ \sum_{i=1}^n \alpha_i \boldsymbol{v}_i \right]$$
$$= \sum_{i=1}^n \alpha_i \exp(At) \, \boldsymbol{v}_i$$



Figure 1: Hyperbolic linear systems in  $\mathbb{R}^3$ .

$$= \sum_{i=1}^{n} \alpha_{i} \left[ \boldsymbol{v}_{i} + tA\boldsymbol{v}_{i} + \frac{1}{2!}t^{2}A^{2}\boldsymbol{v}_{i} + \frac{1}{3!}t^{3}A^{3}\boldsymbol{v}_{i} + \dots \right]$$

$$= \sum_{i=1}^{n} \alpha_{i} \left[ \boldsymbol{v}_{i} + t\lambda_{i}\boldsymbol{v}_{i} + \frac{1}{2!}t^{2}\lambda_{i}^{2}\boldsymbol{v}_{i} + \frac{1}{3!}t^{3}\lambda_{i}^{3}\boldsymbol{v}_{i} + \dots \right]$$

$$= \sum_{i=1}^{n} \alpha_{i} \left[ I + t\lambda_{i} + \frac{1}{2!}t^{2}\lambda_{i}^{2} + \frac{1}{3!}t^{3}\lambda_{i}^{3} + \dots \right] \boldsymbol{v}_{i}$$

$$= \sum_{i=1}^{n} \alpha_{i} \exp\left(\lambda_{i}t\right) \boldsymbol{v}_{i}.$$
(14)

Therefore, we can express the solution of the initial value problem using the eigenvalues/eigenvectors of the matrix A and solving equation (13). Consequently the eigendecomposition of A defines the topology of the general solution of the equation (11).

If  $\lambda_i$  has real part non-null for i = 1, 2, ..., n, we say that the singular point x = 0 is hyperbolic. For n = 3 we have the cases pictured in Figure 1 for hyperbolic systems [3].

Now, consider the vector field  $F : \mathbb{R}^n \to \mathbb{R}^n$ , with F smooth and:

$$\boldsymbol{F}(\boldsymbol{x}) = (F_1(x_1, x_2, \dots, x_n), F_2(x_1, x_2, \dots, x_n), \dots, F_n(x_1, x_2, \dots, x_n)).$$
(15)

Then, we define the Jacobian matrix:

$$D\boldsymbol{F}(\boldsymbol{x}) = \begin{bmatrix} \frac{\partial F_1}{\partial x_1} & \frac{\partial F_1}{\partial x_2} & \cdots & \frac{\partial F_1}{\partial x_n} \\ \frac{\partial F_2}{\partial x_1} & \frac{\partial F_2}{\partial x_2} & \cdots & \frac{\partial F_2}{\partial x_n} \\ \vdots & \vdots & \vdots \\ \frac{\partial F_n}{\partial x_1} & \frac{\partial F_n}{\partial x_2} & \cdots & \frac{\partial F_n}{\partial x_n} \end{bmatrix}.$$

Definition: A hyperbolic singular point of F is a point  $p \in \mathbb{R}^n$  such that: F(p) = 0 and all the eigenvalues of DF(p) have real part non-null.

Hartman's Theorem. Let a smooth vector field  $F : \Re^n \to \Re^n$  and  $p \in \mathbb{R}^n$  a hyperbolic singular point. Then, there are neighbourhoods  $V(p) \subset \mathbb{R}^n$  and  $W(0) \subset \mathbb{R}^n$  such that  $F : V(p) \longrightarrow \mathbb{R}^n$  can be represented by the solutions of:

$$\frac{d\boldsymbol{x}}{dt} = D\boldsymbol{F}\left(\boldsymbol{p}\right)\boldsymbol{x},$$

in the neighbourhood  $W(\mathbf{0})$ .

#### 6 Complex Eigenvalues

For real matrix A and a complex eigenvector  $v = v_1 + iv_2$  with complex eigenvalue  $\lambda = \gamma + i\beta$  we have:

$$A \boldsymbol{v} = \lambda \boldsymbol{v} \implies A \overline{\boldsymbol{v}} = \overline{\lambda} \overline{\boldsymbol{v}}.$$

Therefore,  $\overline{v} = v_1 - iv_2$  is an eigenvector of A with eigenvalue  $\overline{\lambda} = \gamma - i\beta$ . For instance, if  $A \in \mathbb{R}^{2 \times 2}$ , with complex eigenvalues, then from expression (14), the general solution of (11) is:

$$\boldsymbol{x}(t) = \alpha_1 \exp\left[\left(\gamma + i\beta\right)t\right] \left(\boldsymbol{v}_1 + i\boldsymbol{v}_2\right) + \alpha_2 \exp\left[\left(\gamma - i\beta\right)t\right] \left(\boldsymbol{v}_1 - i\boldsymbol{v}_2\right).$$
(16)

Let us define the function:

$$\boldsymbol{\varphi}(t) = \exp\left[\left(\gamma + i\beta\right)t\right]\left(\boldsymbol{v}_1 + i\boldsymbol{v}_2\right) = \exp\left[\gamma t\right]\left(\cos\left(\beta t\right) + i\sin\left(\beta t\right)\right)\left(\boldsymbol{v}_1 + i\boldsymbol{v}_2\right),$$

 $\Leftrightarrow$ 

 $\boldsymbol{\varphi}(t) = \exp\left[\gamma t\right] \left[ \left(\cos\left(\beta t\right) \boldsymbol{v}_1 - \sin\left(\beta t\right) \boldsymbol{v}_2\right) + i\left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) \right].$ 

With this expression we can build another function:

$$\overline{\boldsymbol{\varphi}}(t) = \exp\left[\gamma t\right] \left[ \left(\cos\left(\beta t\right) \boldsymbol{v}_1 - \sin\left(\beta t\right) \boldsymbol{v}_2\right) - i\left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right), \right]$$

So, we can construct the functions

$$\boldsymbol{x}_{1}\left(t\right) = \frac{1}{2}\left(\boldsymbol{\varphi}\left(t\right) + \overline{\boldsymbol{\varphi}}\left(t\right)\right) = Re\left(\boldsymbol{\varphi}\left(t\right)\right),\tag{17}$$

$$\boldsymbol{x}_{2}(t) = \frac{1}{2i} \left( \boldsymbol{\varphi}(t) + \overline{\boldsymbol{\varphi}}(t) \right) = Im \left( \boldsymbol{\varphi}(t) \right).$$
(18)

We can show that:

$$\boldsymbol{x}_{1}(t) = \exp\left[\gamma t\right] \left(\cos\left(\beta t\right) \boldsymbol{v}_{1} - \sin\left(\beta t\right) \boldsymbol{v}_{2}\right),$$

$$\boldsymbol{x}_{2}(t) = \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_{1} + \cos\left(\beta t\right) \boldsymbol{v}_{2}\right).$$

It is possible to verify that  $\{x_1(t), x_2(t)\}\$  is a linerarly independent set of solutions of expression (11). Consequently, in this case, the general solution of the system (11) can be written as:

$$\boldsymbol{x}(t) = c_1 \boldsymbol{x}_1(t) + c_2 \boldsymbol{x}_2(t).$$
<sup>(19)</sup>

If we write  $c_1 = \rho \cos \omega$  and  $c_2 = \rho \sin \omega$ , then we can re-write expression (19) as:

$$\boldsymbol{x}(t) = \rho \cos \omega \exp\left[\gamma t\right] \left(\cos\left(\beta t\right) \boldsymbol{v}_1 - \sin\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_1 + \cos\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin \omega \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin\left(\beta t\right) \exp\left[\gamma t\right] \left(\sin\left(\beta t\right) \boldsymbol{v}_2\right) + \rho \sin\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) + \rho \sin\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) + \rho \sin\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) \exp\left(\beta t\right) + \rho \sin\left(\beta t\right) \exp\left(\beta t\right) \exp\left($$

that renders:

$$\boldsymbol{x}(t) = \rho \exp\left[\gamma t\right] \left(\cos\omega\cos\left(\beta t\right) + \sin\omega\sin\left(\beta t\right)\right) \boldsymbol{v}_{1} +$$

$$\rho \exp\left[\gamma t\right] \left(\sin\omega \cos\left(\beta t\right) - \cos\omega \sin\left(\beta t\right)\right) \boldsymbol{v}_2,$$

which finally gives:

$$\boldsymbol{x}(t) = \rho \exp\left[\gamma t\right] \left(\cos\left(\omega - \beta t\right) \boldsymbol{v}_1 + \sin\left(\omega - \beta t\right) \boldsymbol{v}_2\right).$$
<sup>(20)</sup>

Case 1. Closed trajectories:  $\gamma = 0$ 

$$\boldsymbol{x}(t) = \rho \left( \cos \left( \omega - \beta t \right) \boldsymbol{v}_1 + \sin \left( \omega - \beta t \right) \boldsymbol{v}_2 \right).$$
(21)

Case 2.  $\gamma < 0$ 

$$\lim_{t \longrightarrow +\infty} \boldsymbol{x}(t) = (0,0).$$

Case 3.  $\gamma > 0$ 

$$\lim_{t \longrightarrow +\infty} \boldsymbol{x}(t) = \boldsymbol{\infty}.$$

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